



The Maxwell-Stefan Equations

ChEn 6603

Outline

-  Diffusion in “ideal,” binary systems
 - Particle dynamics
 - Maxwell-Stefan equations
 - Fick’s Law
-  Diffusion in “ideal” multicomponent systems
 - Example: Stefan tube
 - Matrix form of the Maxwell-Stefan equations
 - Fick’s Law for multicomponent systems
 - Reference velocities again

Particle Dynamics

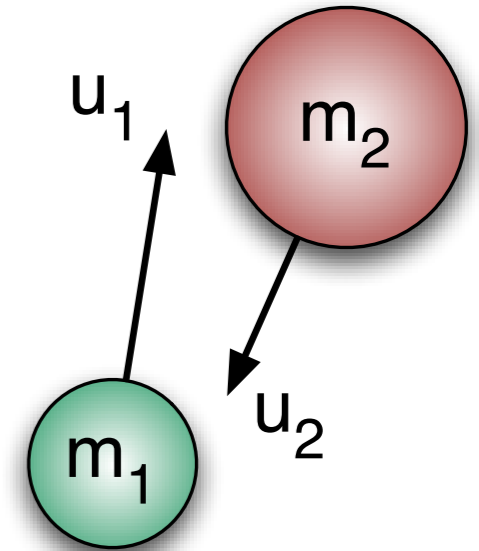
Conservation of momentum:

$$m_1(\mathbf{u}_1 - \mathbf{u}_{f1}) + m_2(\mathbf{u}_2 - \mathbf{u}_{f2}) = 0$$

Conservation of kinetic energy (elastic collision):

$$m_1(\mathbf{u}_1^2 - \mathbf{u}_{f1}^2) + m_2(\mathbf{u}_2^2 - \mathbf{u}_{f2}^2) = 0$$

For molecules, inelastic collisions are known by another name ... what is it?



Solve for final particle velocities:

$$\mathbf{u}_{f1} = \frac{\mathbf{u}_1(m_1 - m_2) + 2m_2\mathbf{u}_2}{m_1 + m_2},$$

$$\mathbf{u}_{f2} = \frac{\mathbf{u}_2(m_2 - m_1) + 2m_1\mathbf{u}_1}{m_1 + m_2}$$

Momentum exchanged in a collision:

$$\begin{aligned} m_1(\mathbf{u}_1 - \mathbf{u}_{f1}) &= m_1\mathbf{u}_1 - \frac{m_1}{m_1 + m_2} (\mathbf{u}_1(m_1 - m_2) + 2m_2\mathbf{u}_2), \\ &= \frac{2m_1m_2(\mathbf{u}_1 - \mathbf{u}_2)}{m_1 + m_2}. \end{aligned}$$

Sum of forces acting on particles of type "1" per unit volume

\propto

Rate of change of momentum of particles of type "1" per unit volume

\propto

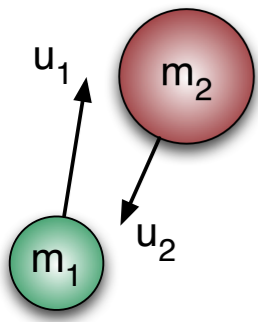
Momentum exchanged per collision between "1" and "2"

\times

Rate of 1-2 collisions per unit volume

$$\mathbf{u}_1 - \mathbf{u}_2$$

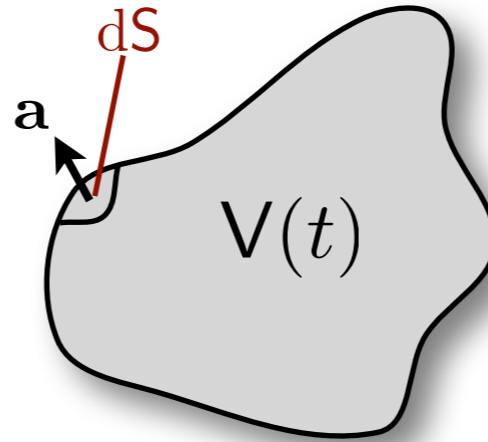
$$\mathcal{X}_1 \mathcal{X}_2$$



Sum of forces acting on particles of type "1" per unit volume \propto Momentum exchanged per collision between "1" and "2" \times Rate of 1-2 collisions per unit volume

Assume:

- System pressure is constant
- Collisions are purely elastic (kinetic energy is conserved in collisions)
- No shear stress (negligible velocity gradients)



$$-\int_{S(t)} x_1 p \mathbf{a} dS = -\int_{V(t)} \nabla(p x_1) dV = -\int_{V(t)} p \nabla x_1 dV$$

Why the (-) sign?

So our force (momentum) balance becomes:

$$-p \nabla x_1 \propto x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2), = f_{12} x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2)$$

f_{12} : drag coefficient for drag that particle "1" feels as a result of interactions with particles of type "2"

Define a "binary diffusion coefficient" as $D_{12} = \frac{p}{f_{12}}$

What is the binary diffusivity a function of?

$$\nabla x_1 = -\frac{x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2)}{D_{12}} \quad \text{Maxwell-Stefan Equations for a binary, ideal mixture.}$$

What about ∇x_2 ?

diffusion driving force for species 1

drag force on species 1 (resisting diffusion)

Fick's Law - Binary Ideal System

$$\nabla x_1 = -\frac{x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2)}{D_{12}}$$

Maxwell-Stefan equations for a binary, ideal system at constant pressure.

$$\nabla x_1 = -\frac{x_2 \mathbf{N}_1 - x_1 \mathbf{N}_2}{c_t D_{12}}$$

$$= -\frac{x_2 \mathbf{J}_1 - x_1 \mathbf{J}_2}{c_t D_{12}}$$

$$\nabla x_1 = -\frac{\mathbf{J}_1}{c_t D_{12}}$$

can you show this?

$$\mathbf{J}_1 = -c_t D_{12} \nabla x_1$$

Fick's law for a binary, ideal system at constant pressure

Re-Cap

- \mathcal{D}_{12} can be interpreted as an inverse drag coefficient.
- $\mathcal{D}_{12} = \mathcal{D}_{21}$ (symmetric due to momentum conservation)
- \mathcal{D}_{12} depends on the characteristics of species 1 and 2 (molecule shapes, etc.), but not on their relative compositions.
- \mathcal{D}_{12} may depend on temperature and pressure.
- We call \mathcal{D}_{12} the “Maxwell-Stefan” diffusivity or “Binary” diffusivity.
- There are no “1-1” interactions here - \mathcal{D}_{11} is not defined.

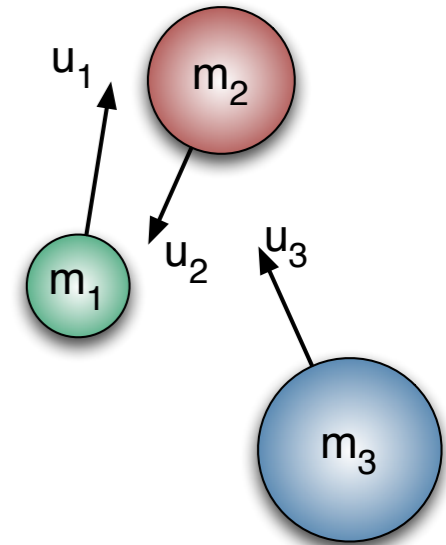
Multicomponent Systems

Binary system: $\nabla x_1 = -\frac{x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2)}{D_{12}}$

Ternary system: must consider 1-2, 1-3, and 2-3 interactions.

$$\nabla x_1 = -\frac{x_1 x_2 (\mathbf{u}_1 - \mathbf{u}_2)}{D_{12}} - \frac{x_1 x_3 (\mathbf{u}_1 - \mathbf{u}_3)}{D_{13}}$$

$$\nabla x_2 = -\frac{x_1 x_2 (\mathbf{u}_2 - \mathbf{u}_1)}{D_{12}} - \frac{x_2 x_3 (\mathbf{u}_2 - \mathbf{u}_3)}{D_{23}}$$



Multicomponent system: must consider i - j interactions.

$$\nabla x_i = -\sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_i x_j (\mathbf{u}_i - \mathbf{u}_j)}{D_{ij}} \quad \xrightarrow{\text{in general...}} \quad \mathbf{d}_i = -\sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_i x_j (\mathbf{u}_i - \mathbf{u}_j)}{D_{ij}}$$

What about $i=j$?

Recall: $\mathbf{N}_i = x_i c \mathbf{u}_i$ $\mathbf{J}_i = \mathbf{N}_i - x_i c \mathbf{u}$

$\mathbf{d}_i = \nabla x_i$
(so far)

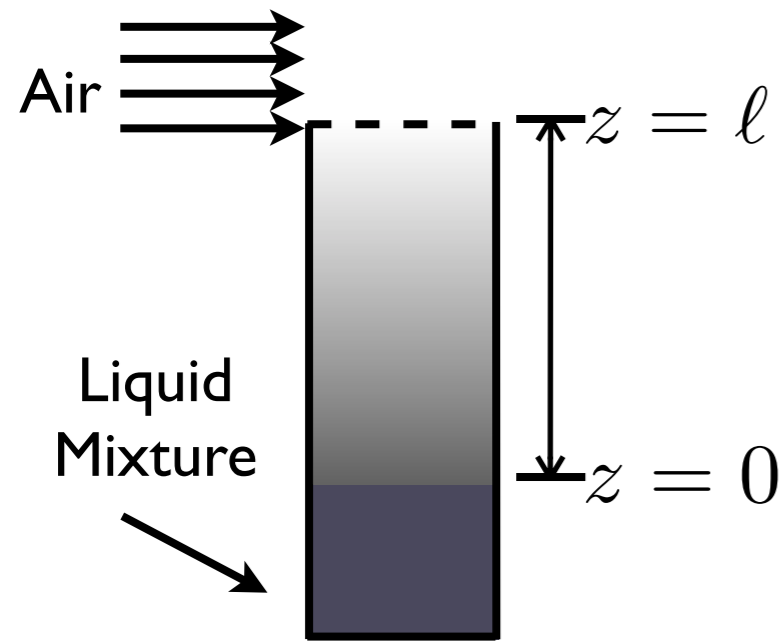
$$\mathbf{d}_i = -\sum_{j \neq i}^n \frac{x_j \mathbf{N}_i - x_i \mathbf{N}_j}{c D_{ij}}$$

$$= -\sum_{j \neq i}^n \frac{x_j \mathbf{J}_i - x_i \mathbf{J}_j}{c D_{ij}}$$

Assumptions:

- Constant Pressure
- Ideal mixture (elastic collisions)
 - ▶ Conservation of translational energy.
 - ▶ Where else could the energy go?

Example: Stefan Tube



Species balance equations (no reaction):

$$\frac{\partial \rho \omega_i}{\partial t} = \frac{\partial \rho_i}{\partial t} = -\nabla \cdot \mathbf{n}_i,$$

$$\frac{\partial c x_i}{\partial t} = \frac{\partial c_i}{\partial t} = -\nabla \cdot \mathbf{N}_i$$

At steady state (1D),

$$\mathbf{n}_i = \alpha_i$$

$$\mathbf{N}_i = \beta_i$$

Convection-
diffusion balance...

Given: D_{ij} , $x_i(z=0)$, $x_i(z=l)$,

$$l = 0.238 \text{ m}, T = 328.5 \text{ K}$$

Find $x_i(z)$

Acetone (1), Methanol (2), Air (3)

$$x_1(z=0) = 0.319, x_2(z=0) = 0.528$$

$$D_{12} = 8.48 \text{ mm}^2/\text{s}$$

$$D_{13} = 13.72 \text{ mm}^2/\text{s}$$

$$D_{23} = 19.91 \text{ mm}^2/\text{s}$$

From the Maxwell-Stefan equations:

$$\frac{dx_i}{dz} = - \sum_{j \neq i}^n \frac{x_j N_i - x_i N_j}{c_t D_{ij}}$$

A semi-analytic solution

Maxwell-Stefan Equations $\frac{dx_i}{dz} = - \sum_{j \neq i}^n \frac{x_j N_i - x_i N_j}{c_t \mathcal{D}_{ij}}$ **Normalized coordinate:** $\eta \equiv \frac{z}{\ell}$, $\frac{d}{dz} = \frac{d}{d\eta} \frac{d\eta}{dz} = \frac{1}{\ell} \frac{d}{d\eta}$

$\frac{1}{\ell} \frac{dx_i}{d\eta} = \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_i N_j - x_j N_i}{c_t \mathcal{D}_{ij}}$, **We need to eliminate x_n from the equation so that we have unknowns $x_1 \dots x_{n-1}$.**

$= x_i \sum_{j \neq i}^n \frac{N_j}{c_t \mathcal{D}_{ij}} - \frac{x_n N_i}{c_t \mathcal{D}_{in}} - \sum_{j \neq i}^{n-1} \frac{x_j N_i}{c_t \mathcal{D}_{ij}}$ **Eliminate x_n by substituting:** $x_n = 1 - \sum_{j=1}^{n-1} x_j = 1 - x_i - \sum_{j \neq i}^{n-1} x_j$

$\frac{1}{\ell} \frac{dx_i}{d\eta} = x_i \sum_{j \neq i}^n \frac{N_j}{c_t \mathcal{D}_{ij}} - \frac{N_i}{c_t \mathcal{D}_{in}} \left(1 - x_i - \sum_{j \neq i}^{n-1} x_j \right) - \sum_{j \neq i}^{n-1} \frac{x_j N_i}{c_t \mathcal{D}_{ij}}$, **rearrange a bit, collecting terms on x_i, x_j .**

$= x_i \left(\frac{N_i}{c_t \mathcal{D}_{in}} + \sum_{j \neq i}^n \frac{N_j}{c_t \mathcal{D}_{ij}} \right) - \frac{N_i}{c_t \mathcal{D}_{in}} + \sum_{j \neq i}^{n-1} \left(\frac{N_i}{c_t \mathcal{D}_{in}} - \frac{N_i}{c_t \mathcal{D}_{ij}} \right) x_j$, **move ℓ over.**

$\frac{dx_i}{d\eta} = \underbrace{\left(\frac{N_i}{c_t \mathcal{D}_{in}/\ell} + \sum_{j \neq i}^n \frac{N_j}{c_t \mathcal{D}_{ij}/\ell} \right)}_{\Phi_{ii}} x_i + \sum_{j \neq i}^{n-1} \underbrace{\left(\frac{N_i}{c_t \mathcal{D}_{in}/\ell} - \frac{N_i}{c_t \mathcal{D}_{ij}/\ell} \right)}_{\Phi_{ij}} x_j - \underbrace{\frac{N_i}{c_t \mathcal{D}_{in}/\ell}}_{\phi_i}$

almost there...

$$\frac{dx_i}{d\eta} = \underbrace{\left(\frac{N_i}{c_t D_{in}/\ell} + \sum_{j \neq i}^n \frac{N_j}{c_t D_{ij}/\ell} \right)}_{\Phi_{ii}} x_i + \sum_{j \neq i}^{n-1} \underbrace{\left(\frac{N_i}{c_t D_{in}/\ell} - \frac{N_i}{c_t D_{ij}/\ell} \right)}_{\Phi_{ij}} x_j - \underbrace{\frac{N_i}{c_t D_{in}/\ell}}_{\phi_i}$$

$$\frac{d(x)}{d\eta} = [\Phi] (x) + (\phi)$$

A system of linear ODEs
with constant coefficients
(c_t, N_j are constant)

$$\Phi_{ii} = \frac{N_i}{c_t D_{in}/\ell} + \sum_{k \neq i}^n \frac{N_k}{c_t D_{ik}/\ell},$$

$$\Phi_{ij} = N_i \left(\frac{1}{c_t D_{in}/\ell} - \frac{1}{c_t D_{ij}/\ell} \right),$$

$$\phi_i = -\frac{N_i}{c_t D_{in}/\ell}$$

Analytic solution
(assuming N_i are all constant)
see T&K §8.3 and Appendix B

$$(x) = \left[\exp [[\Phi] \eta] \right] (x_0) + \left[\exp [[\Phi] \eta] - [I] \right] [\Phi]^{-1} (\phi)$$

Matrix exponential!
 $\exp[\Phi] \neq [\exp(\Phi_{ij})]$!

In Matlab, use "expm"

Note: if we had not eliminated
the " n^{th} " equation, we could not
form the inverses required here.

Algorithm:

1. Guess N_i
2. Calculate $[\Phi], (\phi)$
3. Calculate (x) at $\eta=1$ ($z=\ell$)
4. If (x_ℓ) matches the known boundary condition, we are done. Otherwise return to step 1.

Note: we could also solve the equations numerically in step 3 and eliminate step 2 (work straight from the original Maxwell-Stefan equations)

Matrix Form of Maxwell-Stefan Equations

$$\mathbf{d}_i = - \sum_{j \neq i}^n \frac{x_j \mathbf{N}_i - x_i \mathbf{N}_j}{c \mathcal{D}_{ij}},$$

$$= - \sum_{j \neq i}^n \frac{x_j \mathbf{J}_i - x_i \mathbf{J}_j}{c \mathcal{D}_{ij}}$$

$$\sum_{i=1}^n \mathbf{d}_i = 0$$

Easily shown for the case we have addressed thus far, $\mathbf{d}_i = \nabla x_i$.

Only $n-1$ of these equations are independent.

For a binary system, we have:

$$\nabla x_1 = - \frac{x_2 \mathbf{N}_1 - x_1 \mathbf{N}_2}{c_t \mathcal{D}_{12}},$$

$$\nabla x_2 = - \frac{x_1 \mathbf{N}_2 - x_2 \mathbf{N}_1}{c_t \mathcal{D}_{21}}$$

Show that these sum to zero.

Eliminate \mathbf{J}_n from the set of n equations $\Rightarrow n-1$ equations.

$$\mathbf{J}_n = - \sum_{j=1}^{n-1} \mathbf{J}_j = -\mathbf{J}_i - \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \mathbf{J}_j$$

$$\mathbf{d}_i = - \sum_{j \neq i}^n \frac{x_j \mathbf{J}_i - x_i \mathbf{J}_j}{c_t \mathcal{D}_{ij}},$$

$$c_t \mathbf{d}_i = -\mathbf{J}_i \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{\mathcal{D}_{ij}} + x_i \sum_{\substack{j=1 \\ j \neq i}}^n \frac{\mathbf{J}_j}{\mathcal{D}_{ij}},$$

$$= -\mathbf{J}_i \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{\mathcal{D}_{ij}} + x_i \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \frac{\mathbf{J}_j}{\mathcal{D}_{ij}} + x_i \frac{\mathbf{J}_n}{\mathcal{D}_{in}},$$

Split the summation into individual terms. Recall that we don't have a \mathcal{D}_{ii} term!

Isolate the n^{th} diffusive flux.

$$= -\mathbf{J}_i \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{\mathcal{D}_{ij}} + x_i \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \frac{\mathbf{J}_j}{\mathcal{D}_{ij}} - \frac{x_i}{\mathcal{D}_{in}} \left(\mathbf{J}_i + \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \mathbf{J}_j \right),$$

eliminated the n^{th} diffusive flux

$$= -\mathbf{J}_i \left(\frac{x_i}{\mathcal{D}_{in}} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{\mathcal{D}_{ij}} \right) + x_i \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \left(\frac{1}{\mathcal{D}_{ij}} - \frac{1}{\mathcal{D}_{in}} \right) \mathbf{J}_j$$

Gather \mathbf{J}_i and \mathbf{J}_j terms

$$= -B_{ii} \mathbf{J}_i - \sum_{j \neq i}^{n-1} B_{ij} \mathbf{J}_j$$

Define diagonal and off-diagonal matrix entries.

$$\begin{aligned}
c_t \mathbf{d}_i &= -\mathbf{J}_i \left(\frac{x_i}{D_{in}} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{D_{ij}} \right) + x_i \sum_{\substack{j=1 \\ j \neq i}}^{n-1} \left(\frac{1}{D_{ij}} - \frac{1}{D_{in}} \right) \mathbf{J}_j, \\
&= -B_{ii} \mathbf{J}_i - \sum_{j \neq i}^{n-1} B_{ij} \mathbf{J}_j
\end{aligned}$$

$n-1$ dimensional form
of the Maxwell-Stefan
equations

$n-1$ dimensional
matrix form:

$$c_t(\mathbf{d}) = -[B](\mathbf{J})$$

$$B_{ii} = \frac{x_i}{D_{in}} + \sum_{j \neq i}^n \frac{x_j}{D_{ij}},$$

$$B_{ij} = -x_i \left(\frac{1}{D_{ij}} - \frac{1}{D_{in}} \right)$$

$$c_t \begin{pmatrix} \mathbf{d}_1 \\ \mathbf{d}_2 \\ \vdots \\ \mathbf{d}_{n-1} \end{pmatrix} = - \begin{bmatrix} B_{1,1} & B_{1,2} & \cdots & B_{1,n-1} \\ B_{2,1} & B_{2,2} & \cdots & B_{2,n-1} \\ \vdots & \vdots & \ddots & \vdots \\ B_{n-1,1} & B_{n-1,2} & \cdots & B_{n-1,n-1} \end{bmatrix} \begin{pmatrix} \mathbf{J}_1 \\ \mathbf{J}_2 \\ \vdots \\ \mathbf{J}_{n-1} \end{pmatrix}$$

Note: we can write this in n -dimensional
form, but then $[B]^{-1}$ cannot be formed.

Fick's Law

Maxwell-Stefan Equations
(matrix form)

$$c_t(\mathbf{d}) = -[B](\mathbf{J})$$



Fick's Law
(matrix form)

$$\begin{aligned} (\mathbf{J}) &= -c_t [B]^{-1} (\mathbf{d}) \\ &= -c_t [D] (\nabla x) \end{aligned}$$

so far,
 $\mathbf{d}_i = \nabla x_i$.

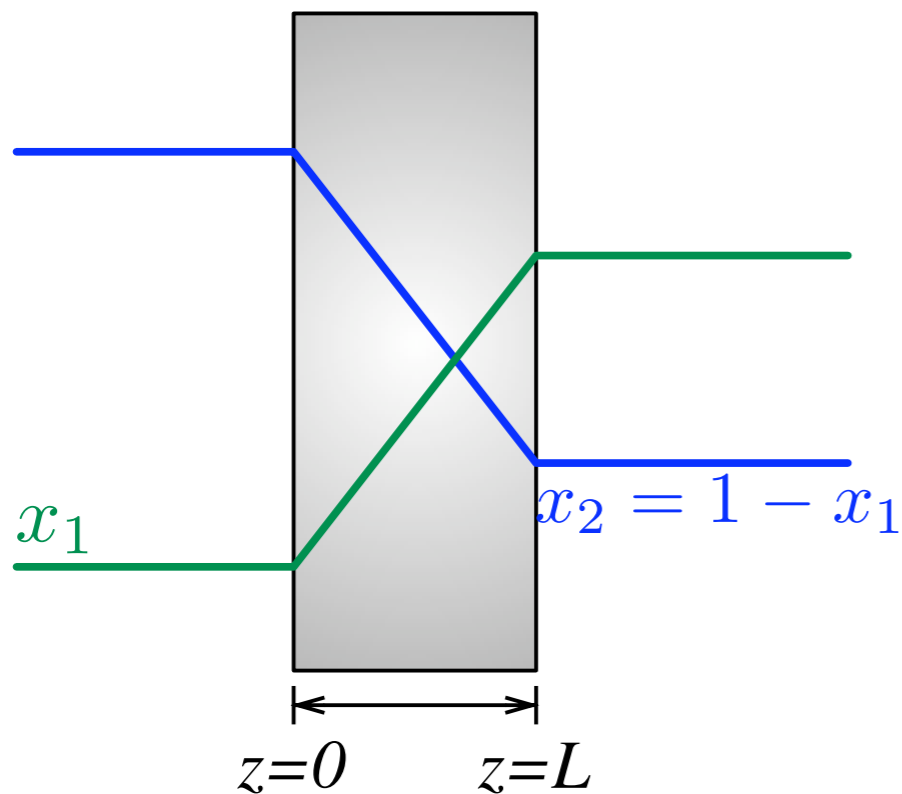
Some Observations:

- For an ideal gas mixture, the \mathcal{D}_{ij} are largely independent of composition (but are functions of T and p), while the D_{ij} are complicated functions of composition.
- The Fickian diffusion coefficients (D_{ij}) may be *negative*, while $\mathcal{D}_{ij} \geq 0$.
- The binary diffusivity matrix is symmetric ($\mathcal{D}_{ij} = \mathcal{D}_{ji}$) but the Fickian diffusivity matrix is not symmetric ($D_{ij} \neq D_{ji}$).
- Note that \mathcal{D}_{ii} never enter in to any expression, and have no physical meaning. However, the Fickian D_{ii} enter directly into the expression for the fluxes, and represent the proportionality constant between the driving force and the diffusion flux for the i^{th} component.
- \mathcal{D}_{ij} are independent of reference frame. D_{ij} is for a molar-averaged velocity reference frame.

Binary/Ternary Comparison

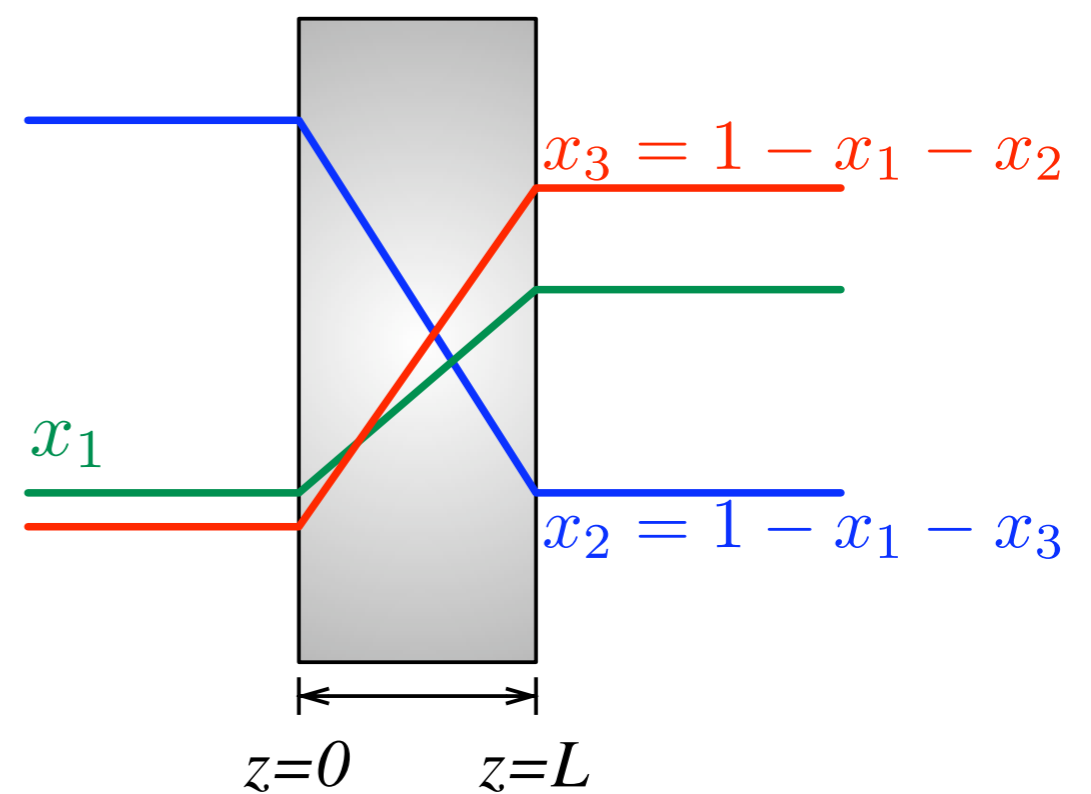
$$\sum_{i=1}^n x_i = 1 \quad \Rightarrow \quad \sum_{i=1}^n \nabla x_i = 0$$

Binary Diffusion



$$\nabla x_2 = -\nabla x_1$$

Ternary Diffusion

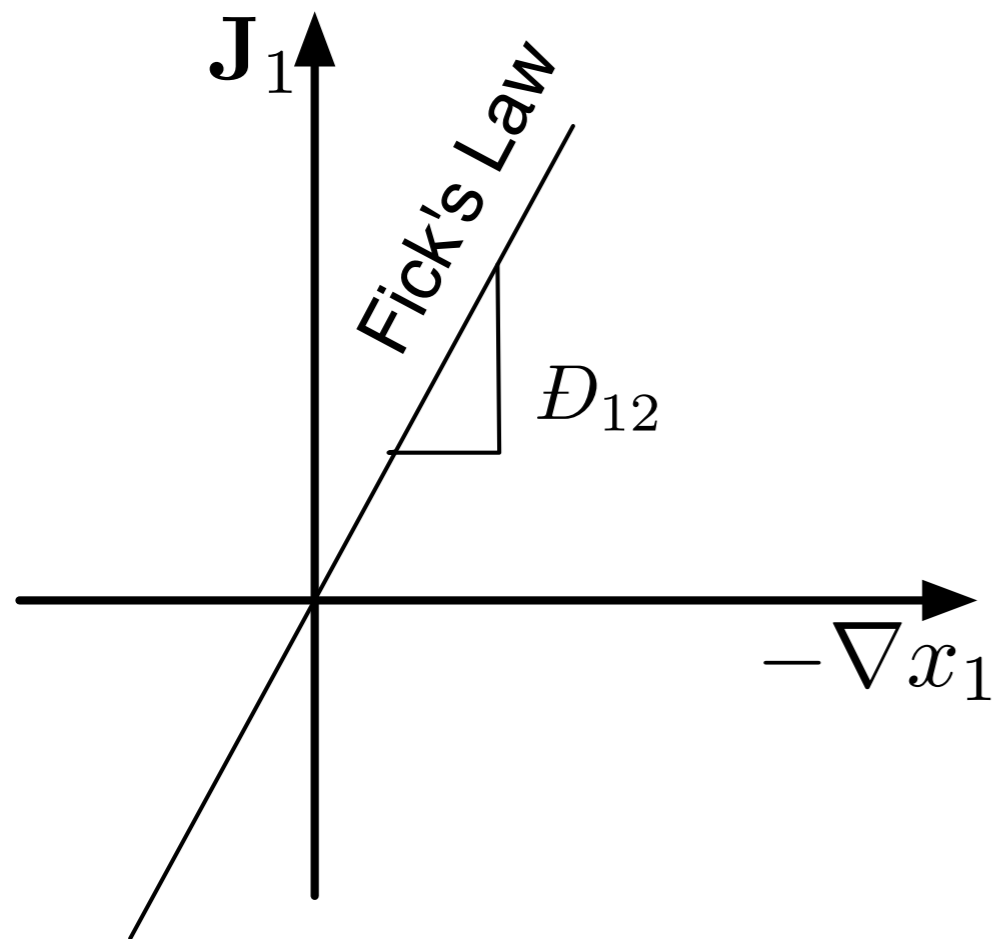


$$\nabla x_1 = -\nabla x_2 - \nabla x_3$$

Diffusion Regimes

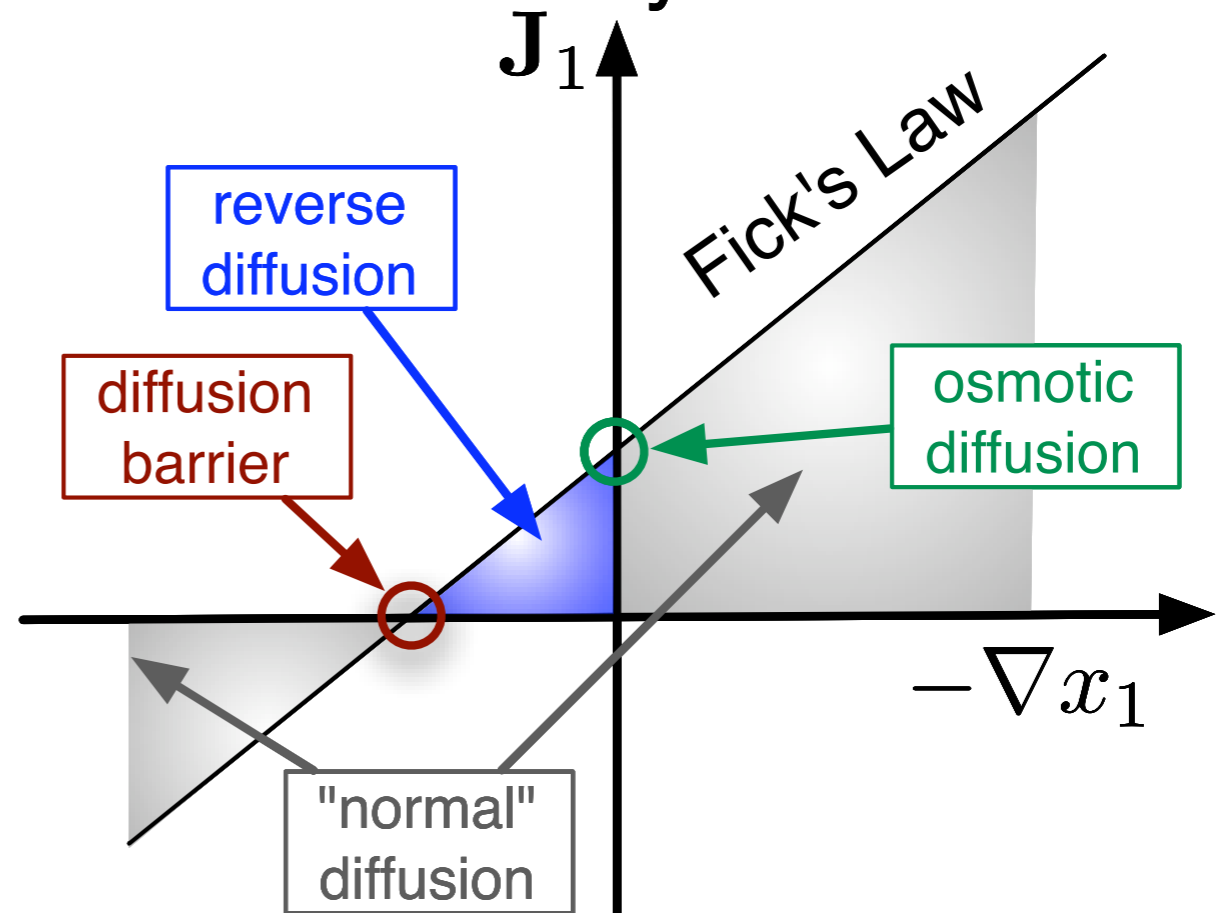
$$(J) = -c_t [D] (\nabla x)$$

Binary Diffusion



$$J_1 = -c_t D \nabla x_1$$

Ternary Diffusion



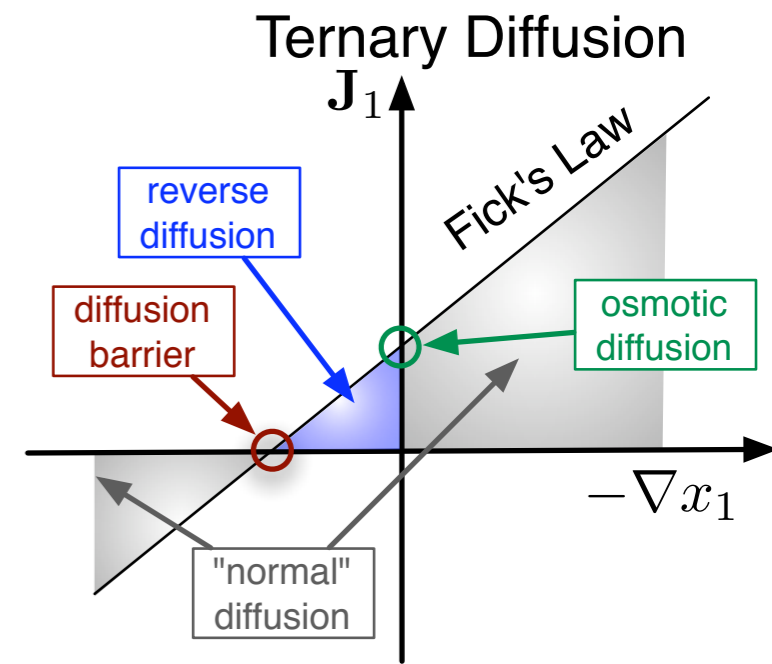
$$J_1 = -c_t D_{11} \nabla x_1 - c_t D_{12} \nabla x_2,$$

$$J_2 = -c_t D_{21} \nabla x_1 - c_t D_{22} \nabla x_2.$$

Multicomponent Effects

$$(\mathbf{J}) = -c_t [B]^{-1} (\mathbf{d})$$

$$\begin{pmatrix} \mathbf{J}_1 \\ \mathbf{J}_2 \\ \vdots \\ \mathbf{J}_{n-1} \end{pmatrix} = -c_t \begin{bmatrix} D_{1,1} & D_{1,2} & \cdots & D_{1,n-1} \\ D_{2,1} & D_{2,2} & \cdots & D_{2,n-1} \\ \vdots & \vdots & \ddots & \vdots \\ D_{n-1,1} & D_{n-1,2} & \cdots & D_{n-1,n-1} \end{bmatrix} \begin{pmatrix} \mathbf{d}_1 \\ \mathbf{d}_2 \\ \vdots \\ \mathbf{d}_{n-1} \end{pmatrix}$$



For multicomponent effects to be important, $D_{ij}\nabla x_j$ must be significant compared to $D_{ii}\nabla x_i$.

$$\left| \frac{D_{ij} \nabla x_j}{D_{ii} \nabla x_i} \right| \sim \mathcal{O}(1)$$

$$|D_{ij}/D_{ii}| \sim \mathcal{O}(1)$$

$$\nabla x_j \neq 0$$

Fick's Law & Reference Velocities

How do we write Fick's law in other reference frames?

$$(\mathbf{J}) = -c[D](\nabla x) \quad \text{Molar diffusive flux relative to a molar-averaged velocity.}$$

$$(\mathbf{j}) = -\rho_t[D^\circ](\nabla\omega) \quad \text{Mass diffusive flux relative to a mass-averaged velocity.}$$

$$(\mathbf{J}^V) = -[D^V](\nabla c) \quad \text{Molar diffusive flux relative to a volume-averaged velocity.}$$

Option 1: Start with GMS equations and write them for the desired diffusive flux and driving force. Then invert to find the appropriate definition for $[D]$.

Option 2: Given $[D]$, define an appropriate transformation to obtain $[D^\circ]$ or $[D^V]$.

$$\begin{aligned} [D^\circ] &= [B^{uo}]^{-1}[\omega][x]^{-1}[D][x][\omega]^{-1}[B^{uo}] & B_{ik}^{uo} &= \delta_{ik} - \omega_i \left(\frac{x_k}{\omega_k} - \frac{x_n}{\omega_n} \right) \\ &= [B^{ou}][\omega][x]^{-1}[D][x][\omega]^{-1}[B^{ou}]^{-1} & B_{ik}^{ou} &= \delta_{ik} - \omega_i \left(1 - \frac{\omega_n x_k}{x_n \omega_k} \right) \end{aligned}$$

$$\begin{aligned} [D^V] &= [B^{Vu}][D][B^{Vu}]^{-1} & B_{ik}^{Vu} &= \delta_{ik} - \frac{x_i}{\bar{V}_t} (\bar{V}_k - \bar{V}_n) \\ &= [B^{Vu}][D][B^{uV}] & B_{ik}^{uV} &= \delta_{ik} - x_i \left(1 - \frac{\bar{V}_k}{\bar{V}_n} \right) \end{aligned}$$

T&K Example 3.2.1

Given $[D^V]$ for the system acetone (1), benzene (2), and methanol (3), calculate $[D]$.

$$(\mathbf{J}) = -c[D](\nabla x) \quad (\mathbf{j}) = -\rho_t[D^\circ](\nabla \omega) \quad (\mathbf{J}^V) = -[D^V](\nabla c)$$

Diffusivities in units of $10^{-9} \text{ m}^2/\text{s}$

$$[D^V] = [B^{Vu}][D][B^{Vu}]^{-1} \\ = [B^{Vu}][D][B^{uV}]$$

$$\bar{V}_1 = 74.1 \times 10^{-6} \frac{\text{m}^3}{\text{mol}}$$

$$\bar{V}_2 = 89.4 \times 10^{-6} \frac{\text{m}^3}{\text{mol}}$$

$$\bar{V}_3 = 40.7 \times 10^{-6} \frac{\text{m}^3}{\text{mol}}$$

x_1	x_2	D_{11}^V	D_{12}^V	D_{21}^V	D_{22}^V
0.350	0.302	3.819	0.420	-0.561	2.133
0.766	0.114	4.440	0.721	-0.834	2.680
0.533	0.790	4.472	0.962	-0.480	2.569
0.400	0.500	4.434	1.866	-0.816	1.668
0.299	0.150	3.192	0.277	-0.191	2.368
0.206	0.548	3.513	0.665	-0.602	1.948
0.102	0.795	3.502	1.204	-1.130	1.124
0.120	0.132	3.115	0.138	-0.227	2.235
0.150	0.298	3.050	0.150	-0.269	2.250